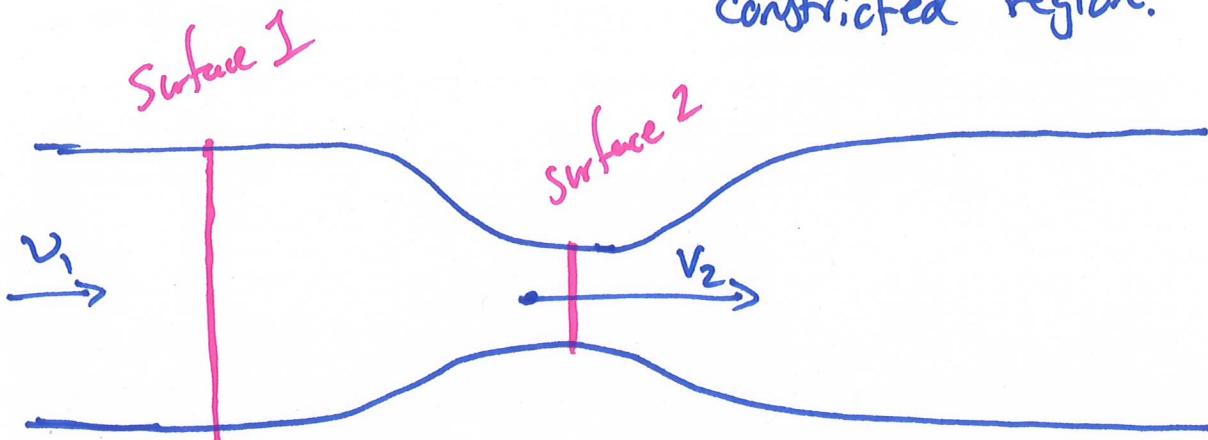


The Venturi tube: a pipe with a constricted region.



Q: How does the velocity change along the center of the pipe?

In the limit of low flow speeds (very subsonic) then the fluid is effectively incompressible.

$$\Delta p \sim \rho \frac{v^2}{2} M^2$$

Mach number $\sim \frac{v}{v_{\text{sound}}}$

density variation from "normal"

if $M \approx 0.3$ then $M^2 \approx 0.1$

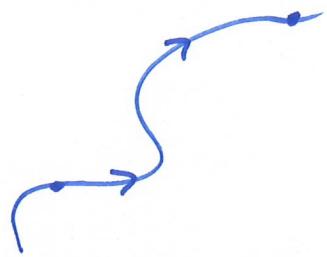
$\Delta p \approx (0.1) \rho$ $\leftarrow \approx 10\%$ change in density

Phy 201 Review:

$$m\vec{a} = \vec{F}_{NET}$$

choose some path in space

integrate along $d\vec{x}$



$$\int m\vec{a} \cdot d\vec{x} = \int \vec{F}_{NET} \cdot d\vec{x}$$



$$d\vec{x} = \vec{v} dt \quad \vec{a} = \frac{d\vec{v}}{dt}$$

$$\int m \left(\frac{d\vec{v}}{dt} \right) \cancel{\vec{v} dt} = W_{NET}$$

$$\int_{V_i}^{V_f} m \vec{v} \cdot d\vec{r}$$

$$\frac{1}{2} m v_f^2 - \frac{1}{2} m v_i^2 = \bullet W_{NET}$$

For low speeds, we can consider constant density along the flow.

Mathematically, this is equivalent to

$$\vec{\nabla} \cdot \vec{v} = 0$$

Conservation of mass told us:

$$\underbrace{\left[\frac{\partial}{\partial t} + \vec{v} \cdot \vec{\nabla} \right]}_{\frac{d}{dt} \rho} \rho = - \rho (\vec{\nabla} \cdot \vec{v})$$

$\frac{d}{dt} \rho \leftrightarrow$ as you follow a small volume of fluid

if $\frac{d}{dt} \rho = 0$, then it must be that $\vec{\nabla} \cdot \vec{v} = 0$.

For low-speed flow, density is constant.

If density is constant then the flux of mass is also constant.

mass flux: $\Gamma = \rho A_{\text{surface}} v$

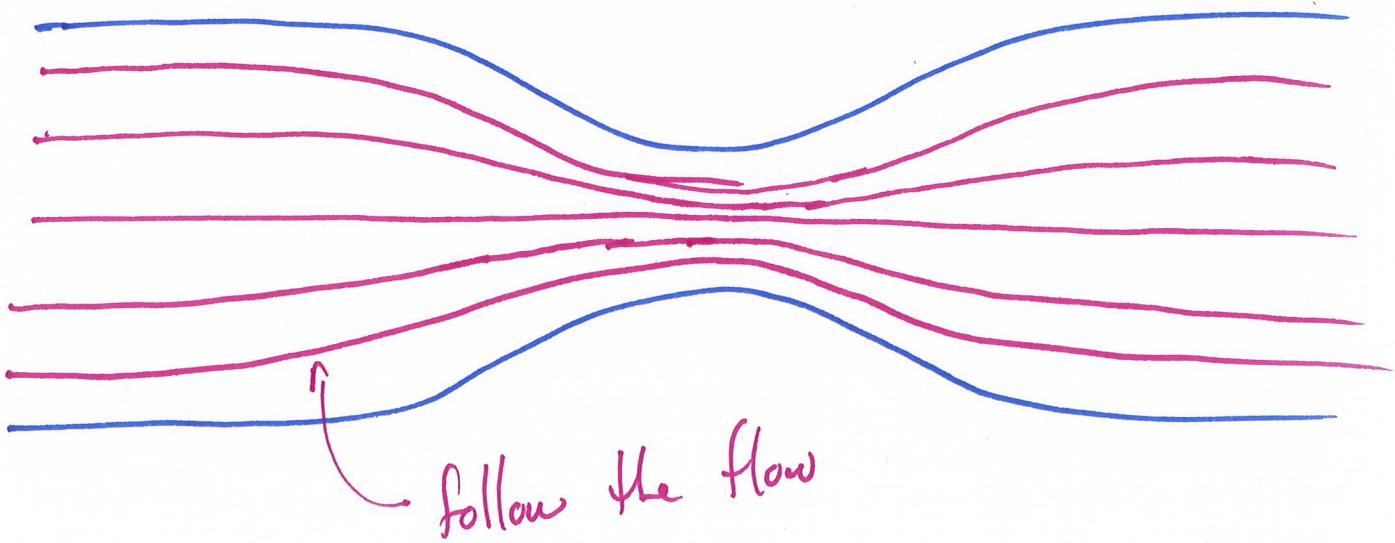
$$\boxed{\text{Surface} \#1} = \boxed{\text{Surface} \#2}$$

$$\cancel{P_1 A_1 V_1} = \cancel{P_2 A_2} \boxed{V_2} \quad P_1 = P_2$$

$$\boxed{V_2 = \frac{A_1}{A_2} V_1}$$

\Rightarrow flows speed up at points of constriction.

A qualitative derivation of Bernoulli's equation.



We now examine N.S. as we follow the flow along a stream line.

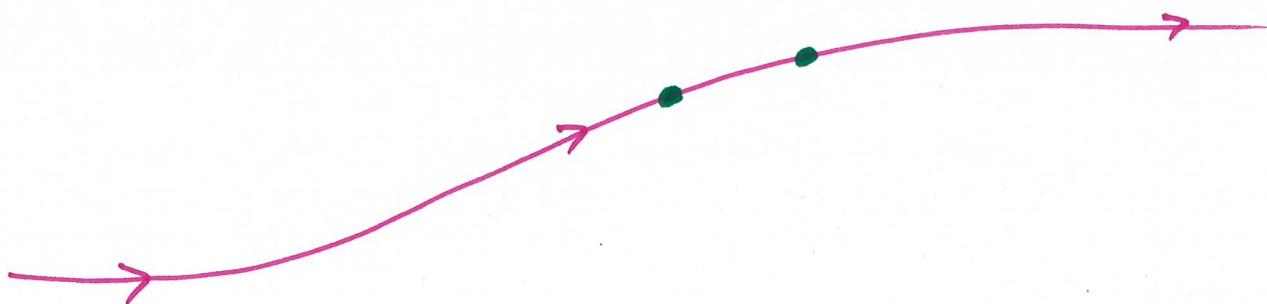
N.S.

$$\cancel{\rho \left[\frac{D}{Dt} + (\vec{v} \cdot \nabla) \vec{v} \right] \vec{v}} = -\vec{\nabla} p + \mu \nabla^2 \vec{v}$$

- look for time-independent solution
"steady-state" flow
- ignore the viscous term

$$\rho (\vec{v} \cdot \nabla) \vec{v} = -\vec{\nabla} p$$

\Rightarrow choose to analyze from a coordinate system aligned with the flow.



along a streamline

$$\rho (\vec{v} \cdot \vec{\nabla}) \vec{v} \Rightarrow \vec{\nabla} \left(\frac{1}{2} \rho v^2 \right)$$

N.S.

$$\vec{\nabla} \left(\frac{1}{2} \rho v^2 \right) = - \vec{\nabla} p$$

$$\vec{\nabla} \left(\frac{1}{2} \rho v^2 + p \right) = 0$$

$$\boxed{\frac{1}{2} \rho v^2 + p = \text{constant}}$$

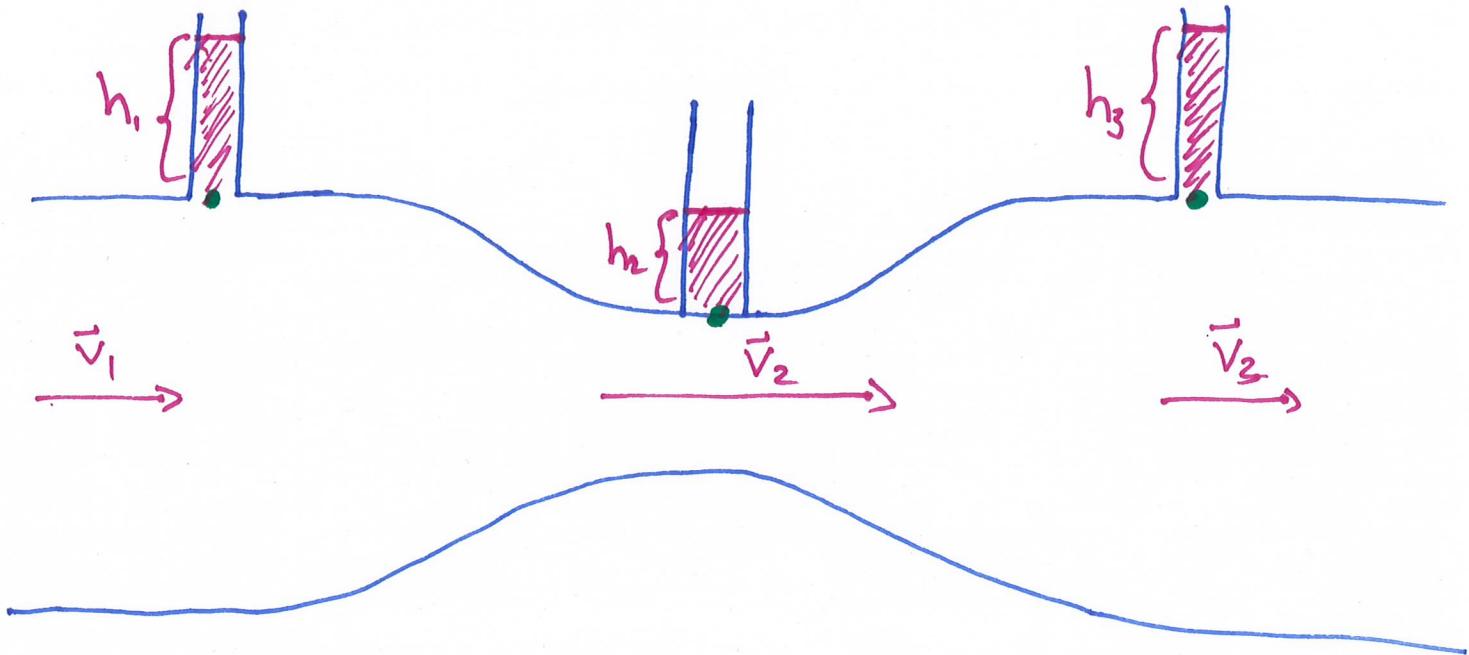
Bernoulli:
Equation (simplified)

To include gravity: add. $- \rho g \hat{z}$

$$- \vec{\nabla}(\rho g z)$$

$$\Rightarrow \frac{1}{2} \rho v^2 + p + \rho g z = \text{constant.}$$

Pressure measurements in a Venturi tube:



At each point along the tube, we have:
 (at bottom of tubes at edge of pipe)

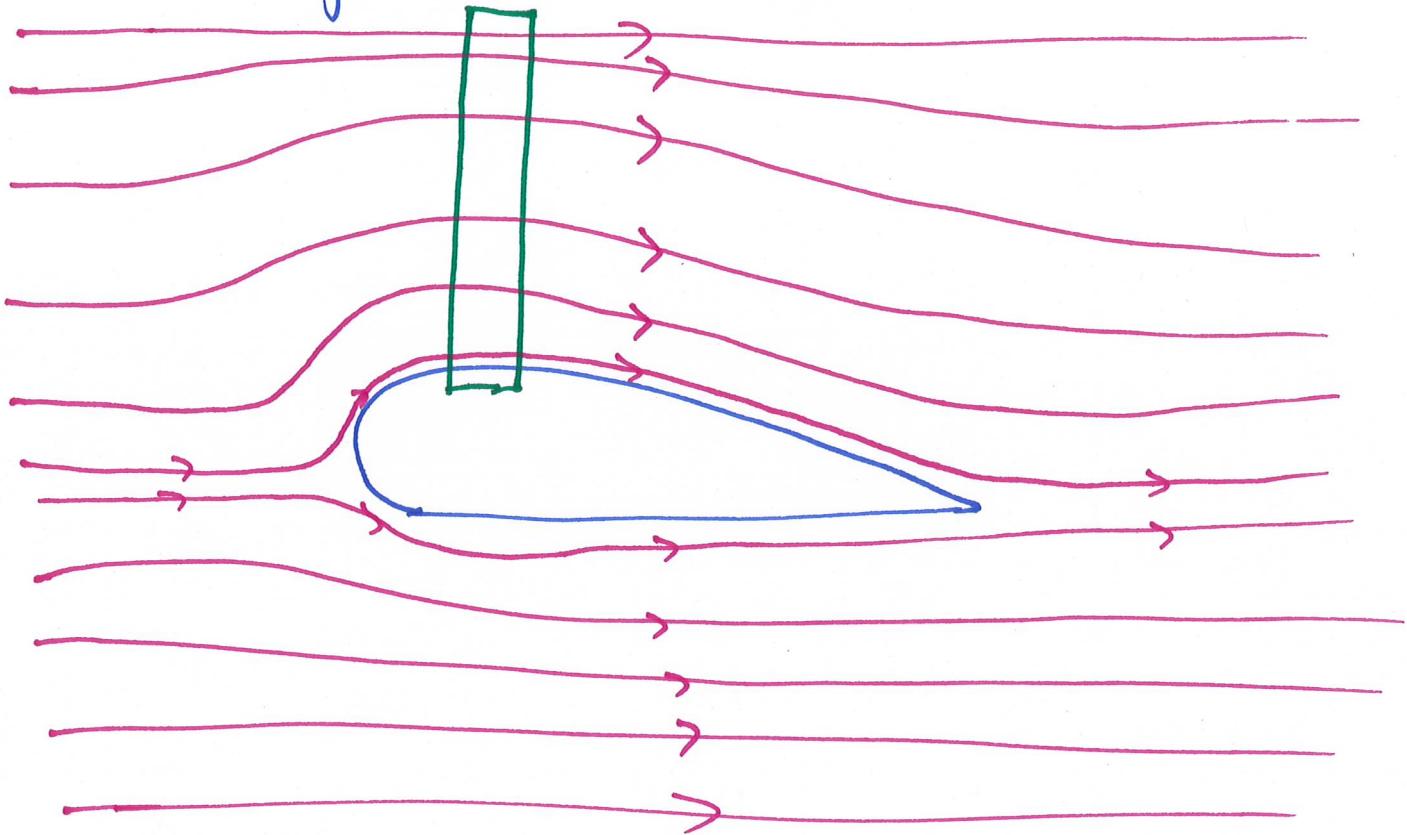
$$\frac{1}{2} \rho v_1^2 + p_1 = \frac{1}{2} \rho v_2^2 + p_2$$

$$p_2 = p_1 + \frac{1}{2} \rho (v_1^2 - v_2^2)$$

$$= p_1 + \frac{1}{2} \rho v_1^2 \left(1 - \underbrace{\left(\frac{A_1}{A_2} \right)^2}_{>1} \right)$$

$$p_2 = p_1 - \frac{1}{2} \rho v_1^2 \left(\frac{A_1^2}{A_2^2} - 1 \right)$$

Lift generation in wings:



Note: far from the wing the flow is undisturbed.

In the region near the wing (green box) the flow looks like that in a Venturi tube.

This compression of the flow means that the flow must accelerate, and by Bernoulli's equation, this flow has less pressure \Rightarrow lift.